

## Surface magnetic states of Ni nanochains modified by using different organic surfactants

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

2010 J. Phys.: Condens. Matter 22 126003

(<http://iopscience.iop.org/0953-8984/22/12/126003>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 129.252.86.83

The article was downloaded on 30/05/2010 at 07:39

Please note that [terms and conditions apply](#).

# Surface magnetic states of Ni nanochains modified by using different organic surfactants

Weimeng Chen<sup>1</sup>, Wei Zhou<sup>2</sup>, Lin He<sup>1</sup>, Chinping Chen<sup>1</sup> and Lin Guo<sup>2</sup>

<sup>1</sup> Department of Physics, Peking University, Beijing 100871, People's Republic of China

<sup>2</sup> School of Chemistry and Environment, Beijing University of Aeronautics and Astronautics, Beijing 100191, People's Republic of China

E-mail: [cpchen@pku.edu.cn](mailto:cpchen@pku.edu.cn) and [guolin@buaa.edu.cn](mailto:guolin@buaa.edu.cn)

Received 2 January 2010

Published 8 March 2010

Online at [stacks.iop.org/JPhysCM/22/126003](http://stacks.iop.org/JPhysCM/22/126003)

## Abstract

Three powder samples of Ni nanochains formed of polycrystalline Ni nanoparticles with an estimated diameter of about 30 nm have been synthesized by a wet chemical method using different organic surfactants. These samples, having magnetically/structurally core-shell structures, all with a ferromagnetic Ni core, are Ni@Ni<sub>3</sub>C nanochains, Ni@Ni<sup>SG</sup> nanochains with a spin glass (SG) surface layer, and Ni@Ni<sup>NM</sup> nanochains with a nonmagnetic (NM) surface layer. The average thickness of the shell for these three samples is determined as about 2 nm. Magnetic properties tailored by the different surface magnetism are studied. In particular, suppression in the saturation magnetization, usually observed with magnetic nanoparticles, is revealed to arise from the surface magnetic states with the present samples.

## 1. Introduction

Surface magnetic states have been a subject of intense study in the past few decades. It has become increasingly important for the emerging field of nanomagnetism. For biomedical applications, magnetic nanoparticles play an important role for the bonding of antibiotics, nucleotides, vitamins, peptides, etc [1–3]. These depend greatly on the magnetic properties of the nanoparticles, especially, the surface magnetism. In information storage, the surface anisotropy plays an important role. It offers an extra degree of freedom to tune the magnetic anisotropy energy, making it attractive for basic investigations [4, 5]. The fundamental role of exchange bias, widely counted on in spin valve and tunneling devices, is directly related to the interaction from the interface or surface [6]. Therefore, the investigations of surface magnetic states are interesting and of great importance.

Theoretical and experimental works reveal that surface atoms may have either enhanced or quenched moments, depending on their chemical environment [7, 8]. Enhancement of the saturation magnetization,  $M_S$ , has been reported in small-sized, elemental metallic clusters [9]. However, many other thin films and nanoparticles demonstrate otherwise [10].

It is reported that the coating of organic molecules [11], CO chemisorption [12] or carbonyl ligation [13] on the surface of magnetic nanoparticles dramatically affects the magnetic properties. The low temperature magnetization reaches only 75% of the bulk value for NiFe<sub>2</sub>O<sub>4</sub> nanoparticles coated with organic molecules [11]. The chemisorption of CO on the metallic Ni surface leads to the quenching of Ni magnetic moments, because electrons of the carbonyl ligation drive the Ni 4s electrons to fill up the 3d shell by repulsive interaction. The Ni atoms in the surface layer, therefore, become nonmagnetic (NM), leaving the magnetism of the inner core unaffected [12, 13]. In many cases, surface magnetic effects are usually featured with a surface spin glass (SG) state and an appreciable reduction of saturation magnetization. For instance, surface SG behavior has often been observed with ferromagnetic (FM) or ferrimagnetic nanoparticles, such as 6.5 nm NiFe<sub>2</sub>O<sub>4</sub> [14], 9–10 nm  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> [15], 12 nm Fe nanoparticles [16] and Ni nanochains [17]. Surface SG properties are even reported with antiferromagnetic (AFM) Co<sub>3</sub>O<sub>4</sub> nanowires having a magnetic core-shell structure of Co<sub>3</sub>O<sub>4</sub><sup>AFM</sup>@Co<sub>3</sub>O<sub>4</sub><sup>SG</sup> [18]. Surface magnetism has also been observed with the partially oxidized composite nanoparticles, Cu@(Cu<sub>2</sub>O<sup>NM</sup> + CuO<sup>AFM</sup>) [19]. Nevertheless, there are

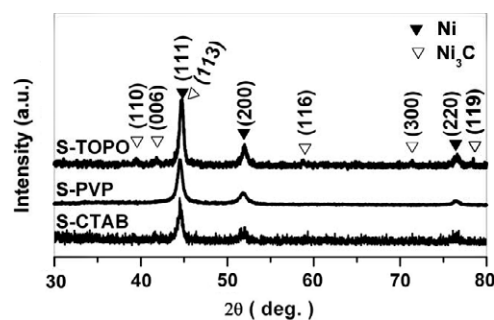
other origins leading to a surface ‘dead layer’ of magnetic nanoparticles and resulting in the suppression of saturation magnetization [20]. These indicate that the surface magnetic state is one of the most important factors dictating the magnetic properties of nanoparticles.

Ni nanochains with a dendritic morphology have been fabricated previously using polyvinyl pyrrolidone (PVP) as a surface modifier [21, 22]. These samples show a SG behavior in addition to a FM phase. The SG behaviors are analyzed from the results of magnetization measurements, and are revealed to come from the surface layer [17]. It has been suggested that the surface modifier (PVP) actually tailors the surface magnetic state. In order to further investigate magnetic properties, such as the finite size effect of the ferromagnetic (FM) transition point with the Ni nanochains, the diameter dependent magnetization reversal behavior with the quasi-one-dimensional chain-like nanostructure, and the possible transport properties with the nanochain structures, it is important to understand the surface magnetic properties and the thickness of the surface layer. We investigate three powder samples of Ni nanochains with a magnetically/structurally core–shell structure, synthesized by a wet chemical method using different surface modifying agents. The estimated outer diameter of the three samples is about 30 nm. They include Ni@Ni<sub>3</sub>C synthesized using trioctylphosphine oxide (TOPO), and Ni nanochains of two kinds synthesized using PVP and hexadecyltrimethyl ammonium bromide (CTAB), forming the magnetic core–shell structure of Ni@Ni<sup>SG</sup> and Ni@Ni<sup>NM</sup>, respectively. They are thus labeled as S-TOPO, S-PVP and S-CTAB. For the three samples, the cores are all of FM Ni, however, with a surface shell layer of different magnetic properties, i.e., a NM surface shell of Ni<sub>3</sub>C for S-TOPO, a magnetically dead layer of Ni (also NM) for S-CTAB, and a SG surface shell for S-PVP.

## 2. Sample preparation and characterization

The detailed processes of synthesis are reported elsewhere. The process and characterization for S-TOPO are reported in [23], S-PVP is Sample B in [22]. The process and the chemical reagents used to prepare S-CTAB are the same as those reported in [24], however, with a different reaction temperature at 197 °C. The chemical reagents used for the preparation of S-CTAB are analytical grade without further purification. In a typical experiment, 0.5 mmol NiCl<sub>2</sub>·6H<sub>2</sub>O and 3 mmol CTAB were dissolved in 60 ml of glycol solvent. Afterward, a solution of hydrazine hydrate (50% v/v), by 1 ml, was dropped into the mixture. When the solution had mixed homogeneously, it was heated to the boiling point (~197 °C), and kept for 5 h. The as-obtained sample was washed with ethanol and deionized water.

The crystal structures were characterized by powder x-ray diffraction (XRD) using a Rigaku Dmax 2200 x-ray diffractometer with Cu K $\alpha$  radiation ( $\lambda = 0.1542$  nm). Transmission electron microscopy (TEM) and high-resolution TEM (HRTEM) investigations were carried using a JEOL JEM-2100F microscope, equipped with EDS (energy dispersive x-ray spectroscopy). The thermogravimetric (TG) analysis



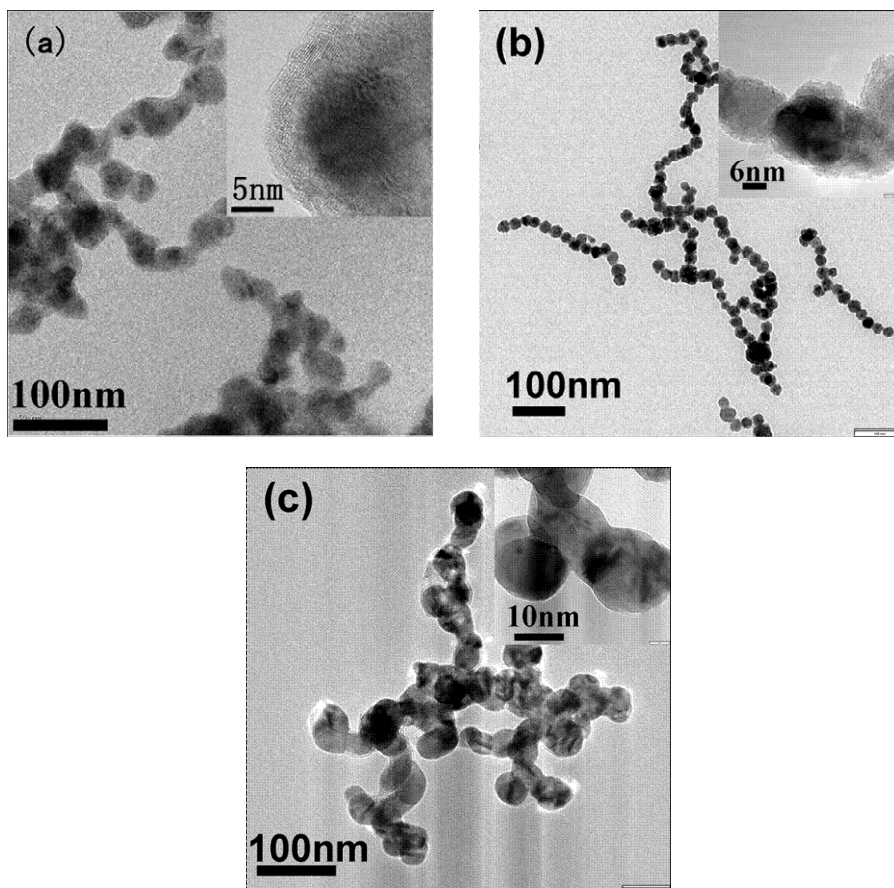
**Figure 1.** XRD patterns for samples modified with TOPO (S-TOPO), PVP (S-PVP) and CTAB (S-CTAB).

was performed using a Diamond Thermogravimetric Analyzer (Perkin-Elmer instruments) under a stream of air. The product was heated from 50 to 600 °C at a rate of 10 °C min<sup>-1</sup>.

Figure 1 shows the XRD patterns for S-TOPO, S-PVP and S-CTAB. It reveals that S-PVP and S-CTAB have a pure Ni phase with a face-centered cubic (fcc) structure. The diffraction peaks correspond to the planes of (111), (200) and (220) of fcc Ni (JCPDS 04-0850). Besides the nickel phase, S-TOPO also contains the Ni<sub>3</sub>C phase, which has been characterized in detail and reported in [23]. The peaks marked with open triangles could be assigned to Ni<sub>3</sub>C (JCPDS 77-0194), corresponding to the planes of (110), (006), (113), (116), (300), and (119) of Ni<sub>3</sub>C. It is worth noting that the (111) peak of Ni overlaps the (113) peak of Ni<sub>3</sub>C.

Figures 2(a)–(c) show the TEM images of S-TOPO, S-PVP, and S-CTAB, respectively. Their morphology is similar, showing chain-like shape with dendritic structure. Their average diameter is estimated at about 30 nm. The corresponding insets show the magnified images. In the inset of figure 2(a), a HRTEM image of Ni@Ni<sub>3</sub>C with a core–shell structure is presented. There is an almost invisible, and very thin, capped layer outside the Ni<sub>3</sub>C shell. It is known that a residual nonmagnetic, organic capped layer, TOPO in this case, is inevitable, even after several times of thorough rinsing. About 10% mass ratio of the organic capped layer is determined by the TG measurement described below. The Ni<sub>3</sub>C shell is estimated from the inset as about 2 to 4 nm in thickness, which is slightly larger than the average thickness of about 2 nm determined by the magnetic measurements. In the inset of figure 2(b) for S-PVP, the HRTEM image reveals that the Ni chains are covered with a vague layer of organic remnants. A layer of almost invisible organic remnants is also observed for S-CTAB, as shown in the inset of figure 2(c). To further confirm the composition of the samples, EDS measurements were conducted, showing the pure Ni element without any other magnetic elements.

The mass ratio of the nonmagnetic organic capped layer is determined by TG measurements, as shown in figure 3. The mass of the as-prepared samples at room temperature is denoted as  $m_0$ . On heating the samples gradually, the mass first decrease slightly due to the burning of the organic remnants. Then, the mass increases dramatically arising from the oxidation of Ni, forming NiO, denoted as  $m(\text{NiO})$ . The mass of Ni is then calculated according to the relative ratio of



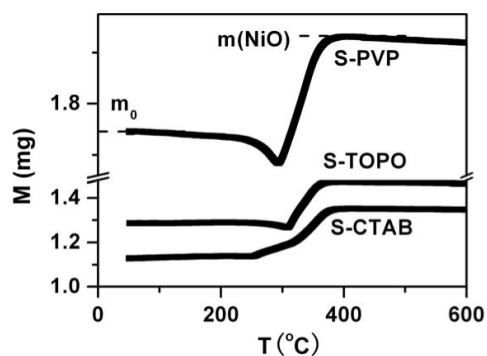
**Figure 2.** Morphology revealed by TEM images for (a) S-TOPO, (b) S-PVP and (c) S-CTAB. The insets are the corresponding HRTEM images.

the formula weight,  $m(\text{Ni}) = (59/75)m(\text{NiO})$ . The correction factor,  $m(\text{Ni})/m_0$ , is then determined as 90%, 87%, and 93% for S-TOPO, S-PVP, and S-CTAB, respectively. The mass correction factor has been accounted for in the normalization of the measured magnetization. It is noted that the mass of C atoms in the  $\text{Ni}_3\text{C}$  shell layer is not taken into account in the above estimation. The NM shell of  $\text{Ni}_3\text{C}$  is treated as if it were a magnetically dead layer of Ni. This introduces about 2% error in mass for  $\text{Ni}_3\text{C}$  being treated as Ni.

### 3. Magnetic measurements and analysis

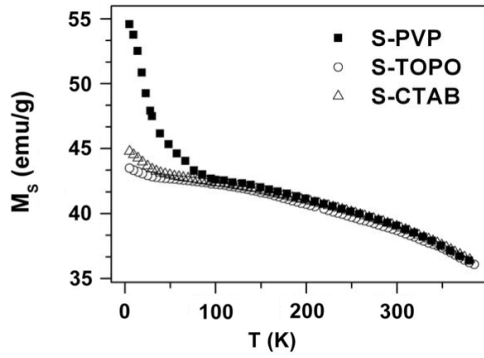
The magnetization measurements were performed on a Quantum Design SQUID magnetometer. The measurements include, (a) temperature dependent saturation magnetization,  $M_S(T)$ , recorded in a field of 20 kOe from  $T = 380$  to 5 K, (b) zero-field-cooling (ZFC) and field-cooling (FC) measurements,  $M_{\text{ZFC}}(T)$  and  $M_{\text{FC}}(T)$  from  $T = 5$  to 380 K, and (c) field dependent hysteresis loops,  $M(H)$ , at a series of fixed temperature from  $T = 5$  to 380 K.

Figure 4 shows the temperature dependence of saturation magnetization,  $M_S(T)$ , of the three samples recorded in an applied field of  $H_{\text{app}} = 20$  kOe. At  $T > 80$  K, the three  $M_S(T)$  curves nearly collapse. The values of  $M_S(T)$  at  $T = 300$  K are determined as  $37.7 \text{ emu g}^{-1}$  (S-TOPO),  $38.2 \text{ emu g}^{-1}$  (S-PVP), and  $38.3 \text{ emu g}^{-1}$  (S-CTAB). They account for about



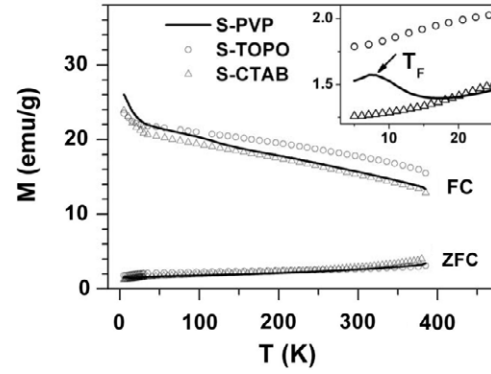
**Figure 3.** Thermogravimetric analyses for the samples of S-TOPO, S-PVP and S-CTAB. The mass of the as-prepared sample is denoted as  $m_0$ , and the final NiO is denoted as  $m_{\text{NiO}}$ . The calculated Ni contents,  $m_{\text{Ni}}$ , are 93%, 90% and 87% for TOPO, CTAB and PVP, respectively.

70% of the corresponding bulk value,  $\sim 54.2 \text{ emu g}^{-1}$ , at 300 K [25]. The reduction of the saturation magnetization is attributed to the magnetically inert property of the surface shell. It is noted that the mass effect of the NM organic capped layer has already been corrected for by the TG measurements. Otherwise, the saturation magnetization per unit mass would have been further reduced by 10–15%. For S-TOPO, the



**Figure 4.** Saturation magnetization,  $M_S(T)$ , recorded in an applied field of 20 kOe from 380 to 5 K. The magnetization of S-PVP increases dramatically at  $T < 80$  K.

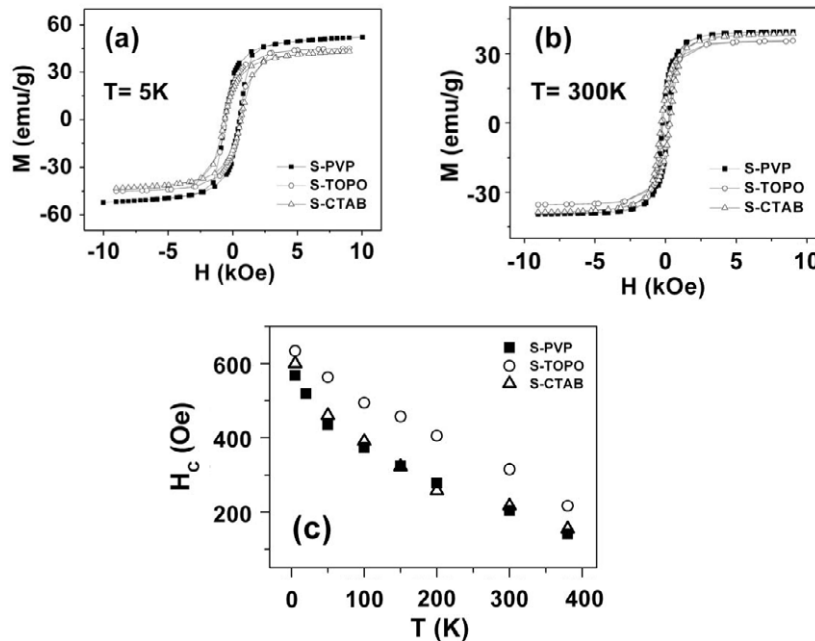
reduction in the saturation magnetization is obvious because the  $\text{Ni}_3\text{C}$  shell is NM [26]. The average diameter of the FM cores is then estimated as  $D_{\text{core}} = 26.6$  nm with a  $\text{Ni}_3\text{C}$  shell thickness of about 1.7 nm. It is consistent with the result of shell thickness, 1–4 nm, observed by the HRTEM investigation [23]. For S-PVP and S-CTAB, although there is no obvious shell structure like the  $\text{Ni}_3\text{C}$  shell of S-TOPO, a FM core of Ni with roughly the same diameter,  $D_{\text{core}} \sim 26$  nm, is also concluded. At  $T < 80$  K,  $M_S(T)$  for S-PVP increases dramatically, reaching the bulk value of Ni,  $\sim 55$  emu  $\text{g}^{-1}$ . It is attributed to the contribution from the surface SG shell [17]. On the other hand, the shell layer of S-CTAB, although of pure Ni, does not show any magnetism at all, very much like S-TOPO with a NM  $\text{Ni}_3\text{C}$  shell layer. Therefore, S-CTAB exhibits properties of a magnetically core-shell structure,  $\text{Ni}@\text{Ni}^{\text{NM}}$ , with a magnetically ‘dead shell’ of  $\text{Ni}^{\text{NM}}$ . The presence of a magnetically ‘dead layer’ had been



**Figure 5.** ZFC and FC  $M(T)$  curves for S-TOPO, S-PVP and S-CTAB measured in an applied field of 90 Oe from 5 to 380 K. The inset shows the ZFC curves in the low temperature region. The peak around 8 K in the ZFC curve with S-PVP is attributed to the freezing of the surface SG shell.

reported in the early days with the surface of Ni films [27]. In addition, the magnetic moments of the surface Ni are reported to be quenched, e.g., by the carbonyl ligation on the surface [13].

Figure 5 shows the  $M_{\text{ZFC}}(T)$  and  $M_{\text{FC}}(T)$  curves. The ZFC curves were recorded in an applied field,  $H_{\text{app}} = 90$  Oe, in the warming process after the sample was cooled down to  $T = 5$  K under zero applied field. For the FC measurement, the procedure of data collection was the same, except that the sample was cooled down to 5 K in an applied magnetic field of 20 kOe. For all of the three samples, the  $M_{\text{FC}}(T)$  and  $M_{\text{ZFC}}(T)$  curves separate from each other as the temperature goes up to 380 K. It indicates that the blocking temperature is higher than 380 K. The inset shows the blown-up  $M_{\text{ZFC}}(T)$  curves in the low temperature region. A freezing peak appears at



**Figure 6.**  $M(H)$  curves for S-TOPO, S-PVP and S-CTAB at (a)  $T = 5$  K, (b)  $T = 300$  K, and (c) coercivity determined from the  $M(H)$  curves at different temperatures.

$T_F \sim 8$  K with S-PVP, similar to those at about 13 K observed for 50, 75 and 150 nm Ni nanochains also synthesized using the surfactant of PVP [17]. It further confirms that S-PVP exhibits the magnetic core-shell structure of Ni@Ni<sup>SG</sup>. For the other two samples, S-TOPO and S-CTAB, there is no such characteristic feature, as is expected for the samples with a NM shell layer.

The  $M(H)$  measurements for the three samples are performed at various temperature from 5 to 380 K. Figure 6(a) shows the  $M(H)$  curves measured at  $T = 5$  K. The magnetization determined in the high field region at  $H = 10$  kOe, and the coercivity,  $H_C$ , are 45.0 emu g<sup>-1</sup> and 600 Oe for S-TOPO, 52.3 emu g<sup>-1</sup> and 568 Oe for S-PVP, and 43.1 emu g<sup>-1</sup> and 634 Oe for S-CTAB, respectively. The magnetization of S-PVP is higher than that of S-TOPO, or S-CTAB by about 20%, attributed to the contribution from the SG shell. The  $M(H)$  curves at  $T = 300$  K are shown in figure 6(b). The difference in saturation magnetization is more or less reduced. The coercivity determined from the  $M(H)$  curves at various temperatures is shown in figure 6(c). The magnetization reversal can be described by the fanning mode based on the chain of sphere model proposed by Jacobs and Bean [28], as discussed in our previous work for S-TOPO [23].

#### 4. Conclusion

We have investigated the magnetic properties of three samples of Ni nanochains, with an estimated outer diameter of 30 nm, synthesized by different surface modifying agents, including CTAB, TOPO and PVP. The surface magnetic properties are modified significantly and differently by the surfactants in the synthesis process, forming different magnetic shell structures. The CTAB leads to nanochains of Ni@Ni<sup>NM</sup> with a magnetically dead layer of Ni surface shell, while the TOPO produces final products of Ni@Ni<sub>3</sub>C nanochains with a shell of NM Ni<sub>3</sub>C. With PVP, a SG shell layer of Ni is formed, resulting in the magnetic structure of Ni@Ni<sup>SG</sup>. The saturation magnetization of these samples accounts for only 70% of the bulk value at 300 K, attributed to the presence of the magnetically inert layer. This value would be further reduced, by more than 10%, if the mass of the outermost capped layer of organic remnants was not corrected for. The thickness of the shell layer is determined as about 2 nm for all of the three samples.

#### Acknowledgments

This work is supported by the NSFC (Grant nos 10874006, 20973019 and 50725208), the State Key Project of Fundamental Research for Nanoscience and Nanotechnology (2006CB932301), the National Basic Research Program of China (nos 2009CB939901, 2010CB934700 and 2010CB934601).

#### References

- [1] Pankhurst Q A, Connolly J, Jones S K and Dobson J 2003 *J. Phys. D: Appl. Phys.* **36** R167
- [2] Tartaj P, Morales M P, Veintemillas-Verdaguer S and Gonzalez-Carreño T 2003 *J. Phys. D: Appl. Phys.* **36** R182
- [3] Lu A, Salabas E L and Schuth F 2007 *Angew. Chem. Int. Edn* **46** 1222
- [4] Gambardella P, Rusponi S, Veronese M, Dhési S S, Grazioli C, Dallmeyer A, Cabria I, Zeller R, Dederichs P H, Kern K, Carbone C and Brune H 2003 *Science* **300** 1130
- [5] Kaneyoshi T 1991 *J. Phys.: Condens. Matter* **3** 4497
- [6] Nogues J, Sort J, Langlais V, Skumryev V, Surinach S, Munoz J S and Baro M D 2005 *Phys. Rep.* **422** 65
- [7] Bader S D 2002 *Surf. Sci.* **500** 172
- [8] Liu F, Press M R, Khanna S N and Jena P 1989 *Phys. Rev. B* **39** 6914
- [9] Jensen P J and Bennemann K H 1995 *Z. Phys. D* **35** 273
- [10] Dormann J L, Fiorani D and Tronc E 1997 *Adv. Chem. Phys.* **98** 283
- [11] Berkowitz A E, Lahut J A, Jacobs I S, Levinson L M and Forester D W 1975 *Phys. Rev. Lett.* **34** 594
- [12] Feigerle C S, Seiler A, Pena J L, Celotta R J and Pierce D T 1986 *Phys. Rev. Lett.* **56** 2207
- [13] van Leeuwen D A, van Ruitenbeek J M, de Jongh L J, Ceriotti A, Pacchioni G, Haberlen O D and Rosch N 1994 *Phys. Rev. Lett.* **73** 1432
- [14] Kodama R H, Berkowitz A E, McNiff E J Jr and Foner S 1996 *Phys. Rev. Lett.* **77** 394
- [15] Martinez B, Obradors X, Balcells L I, Rouanet A and Monty C 1997 *Phys. Rev. Lett.* **80** 181
- [16] Bonetti E, Bianco L D, Fiorani D, Rinaldi D, Caciuffo R and Hernando A 1999 *Phys. Rev. Lett.* **83** 2829
- [17] He L, Zheng W Z, Zhou W, Du H L, Chen C P and Guo L 2007 *J. Phys.: Condens. Matter* **19** 036216
- [18] Benitez M J, Petravic O, Salabas E L, Radu E, Tuysuz H, Schuth F and Zabel H 2008 *Phys. Rev. Lett.* **101** 097206
- [19] Li Q, Zhang S W, Zhang Y and Chen C P 2006 *Nanotechnology* **17** 4981
- [20] Westman C, Jang S, Kim C, He S, Harmon G, Miller N, Graves B, Poudyal N, Sabirianov R, Zeng H, DeMarco M and Liu J P 2008 *J. Phys. D: Appl. Phys.* **41** 225003
- [21] Liu C M, Guo L, Wang R M, Deng Y, Xu H B and Yang S H 2004 *Chem. Commun.* **23** 2726
- [22] Zhou W, He L, Cheng R, Guo L, Chen C P and Wang J L 2009 *J. Phys. Chem. C* **113** 17355
- [23] Zhou W, Zheng K, He L, Wang R M, Guo L, Chen C P, Han X D and Zhang Z 2008 *Nano Lett.* **8** 1147
- [24] Zhou W, Guo L, He L and Chen C P 2008 *Phys. Status Solidi a* **205** 1109
- [25] Sellmyer D J, Zheng M and Skomski R 2001 *J. Phys.: Condens. Matter* **13** R433
- [26] Yue L P, Sabirianov R, Kirkpatrick E M and Leslie-Pelecky D L 2000 *Phys. Rev. B* **62** 8969
- [27] Liebermann L, Clinton J, Edwards D M and Mathon J 1970 *Phys. Rev. Lett.* **25** 232
- [28] Jacobs L S and Bean C P 1955 *Phys. Rev.* **100** 1060